

Introduction to *ab-initio* methods and relation to nanomaterials research

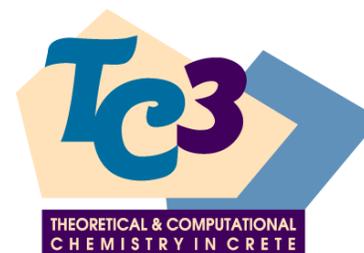


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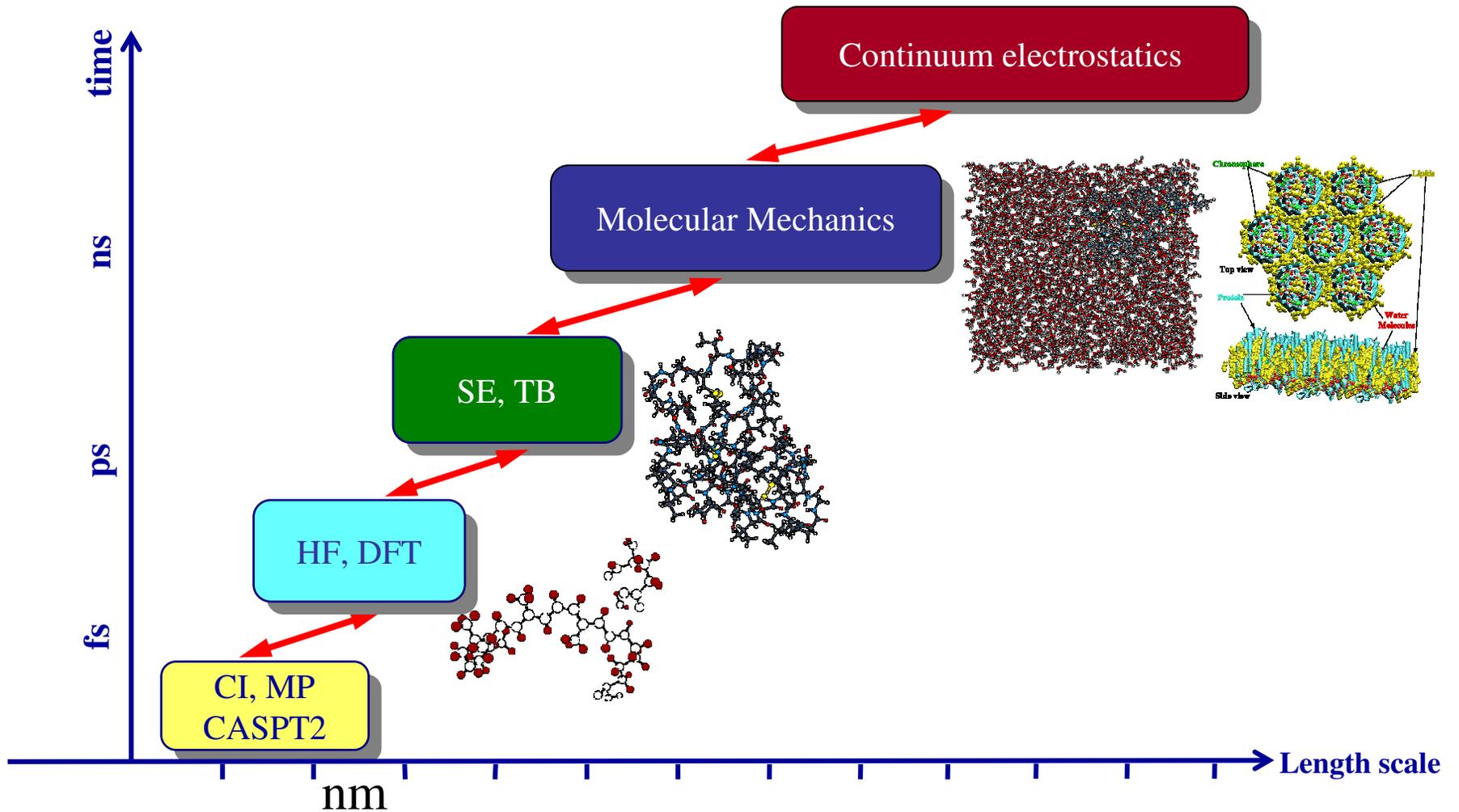
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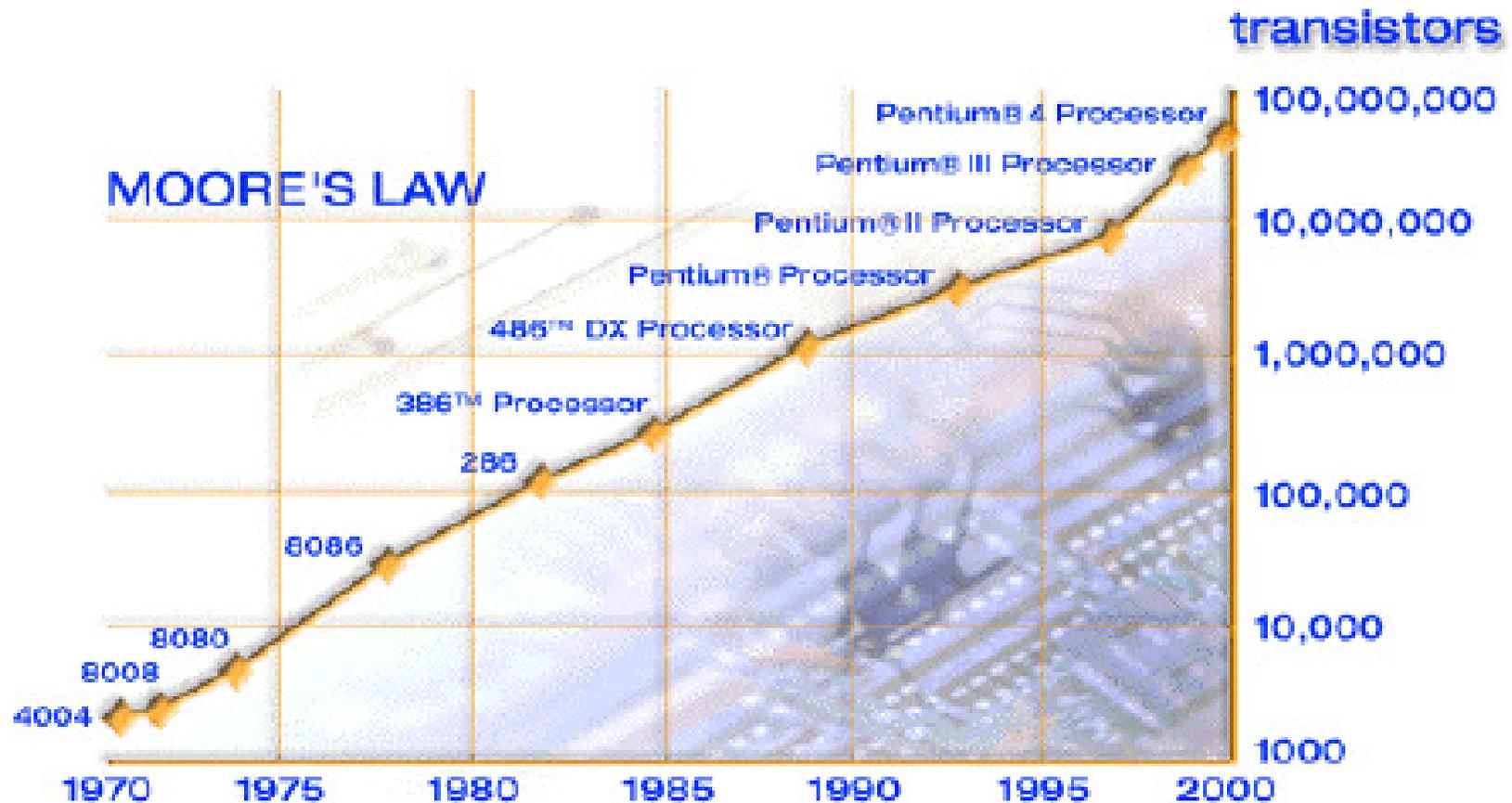
Why we need theory and modeling ?

-  Explanation of experimental results
-  Material testing *in silico*
-  Designing novel materials

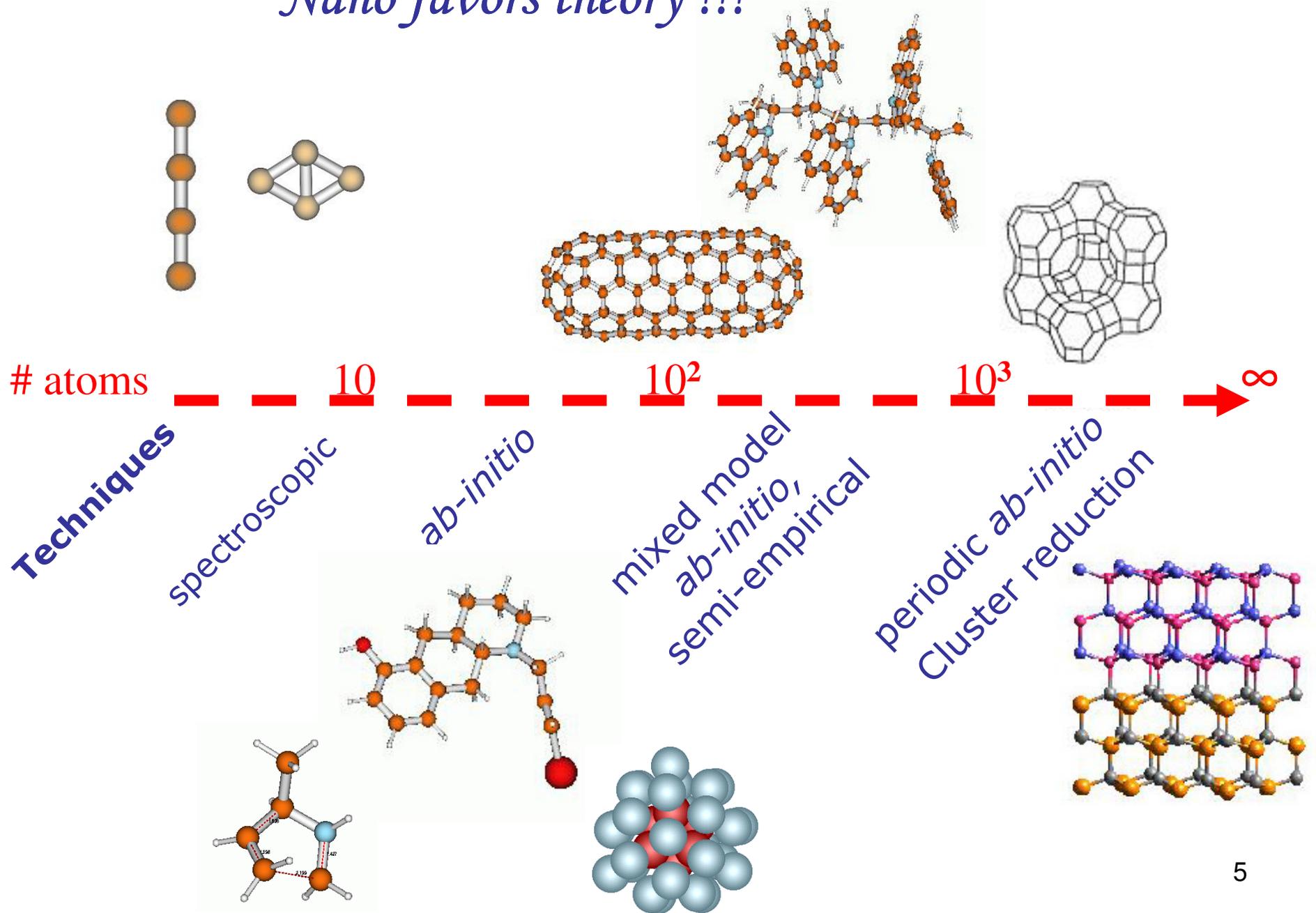
Hierarchy of methods



In past 30 years, computational power
(driven by Moore's Law) has increased
by over 5 orders of magnitude



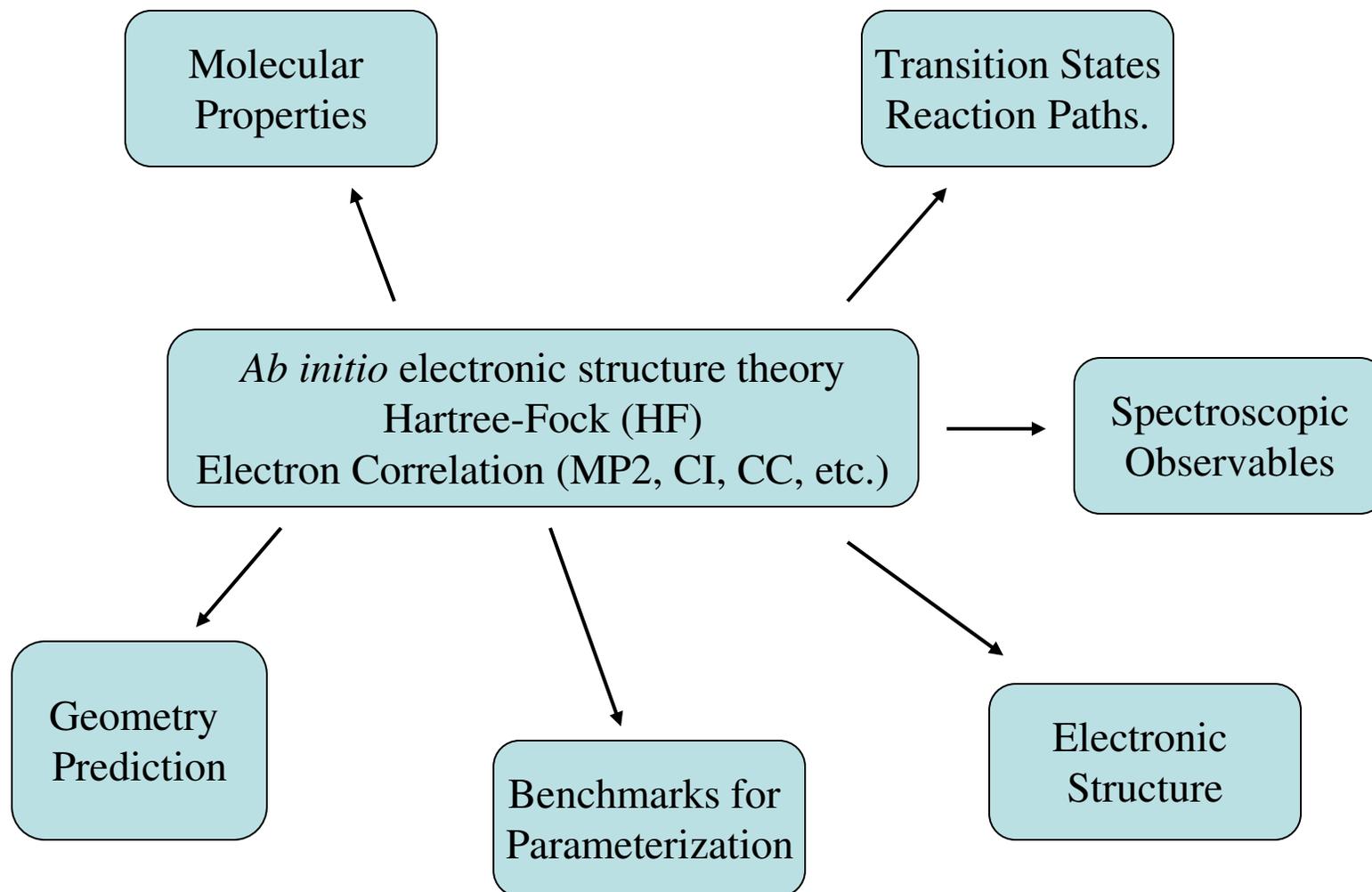
Nano favors theory !!!



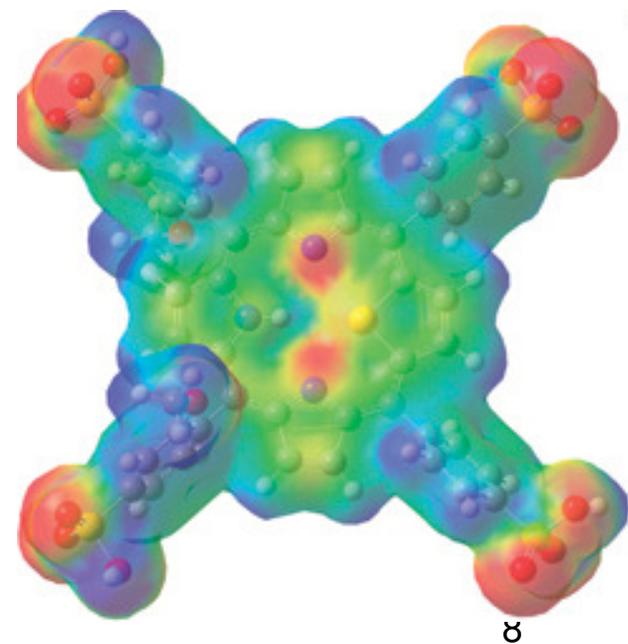
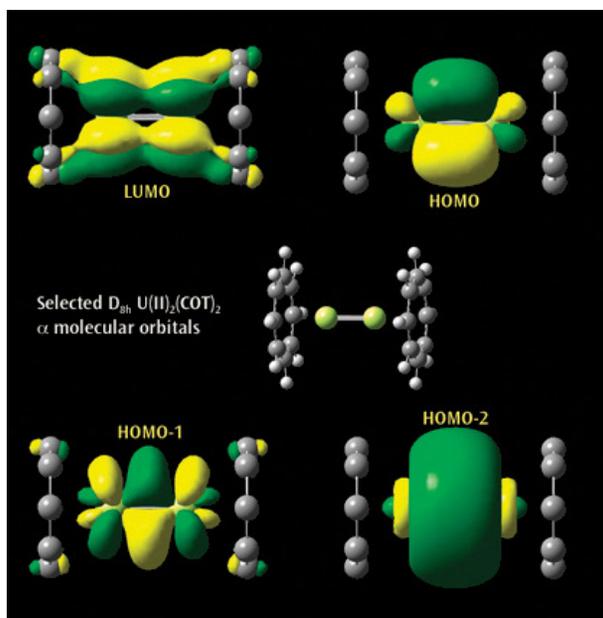
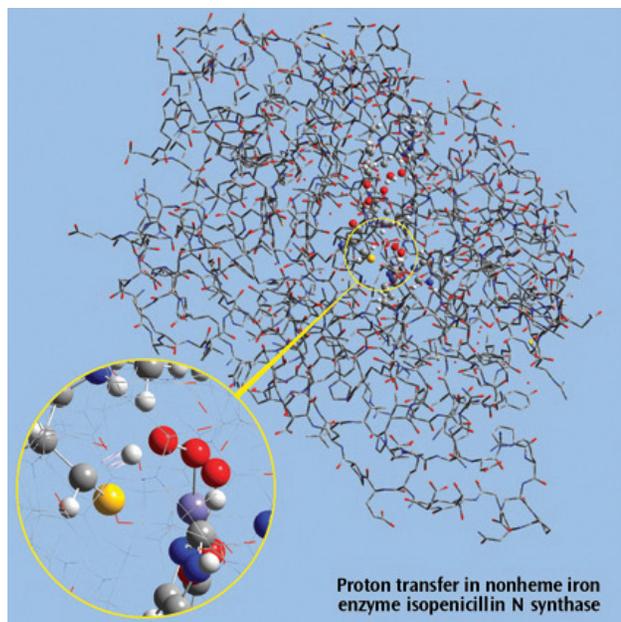
What is Quantum Chemistry?

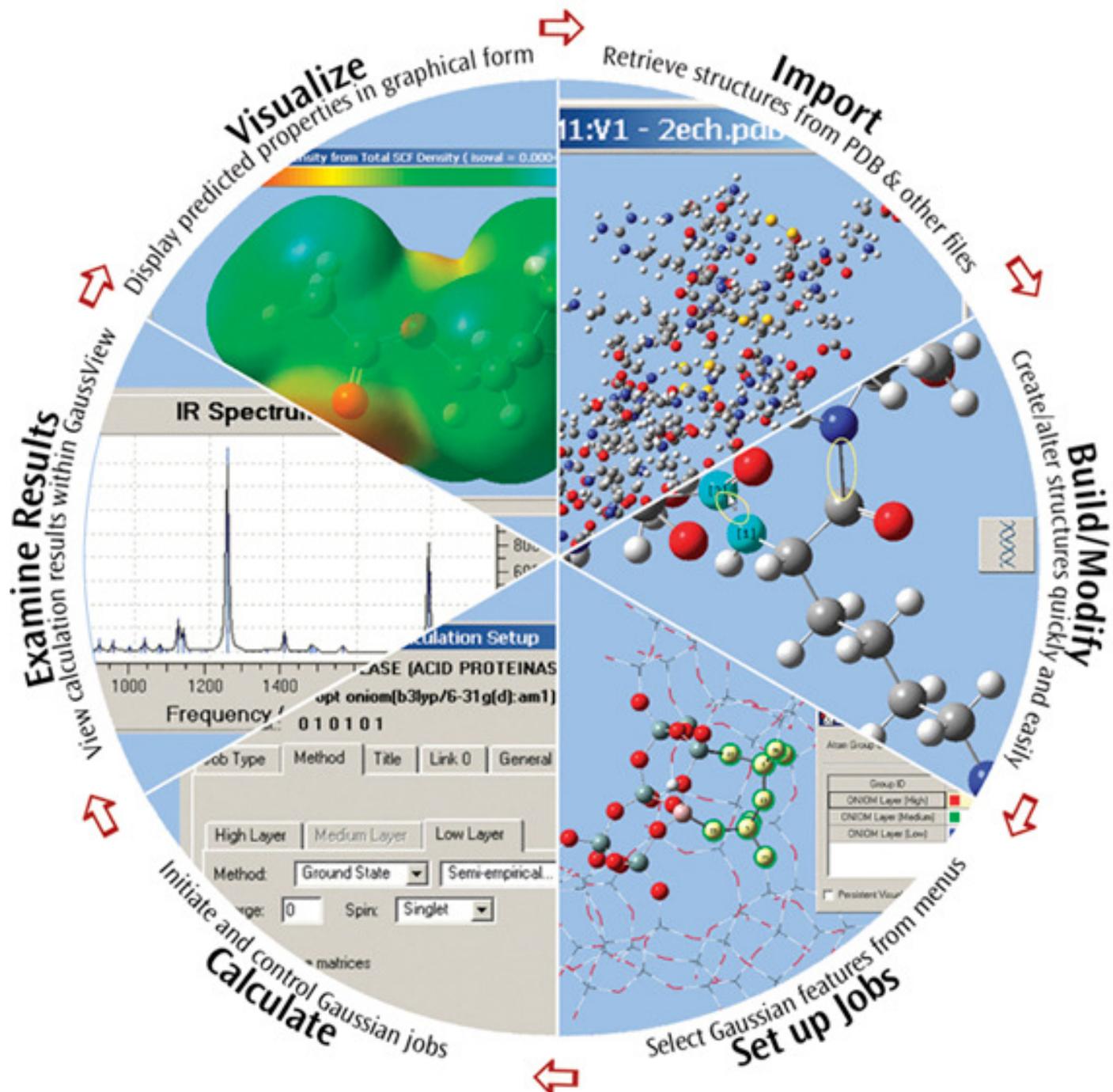


- Quantum Mechanics applied to Atoms and Molecules
- Aim: Understanding of Electronic Structure
- Solution of the electronic Schrödinger equation
- Derived: Properties of Atoms and Molecules



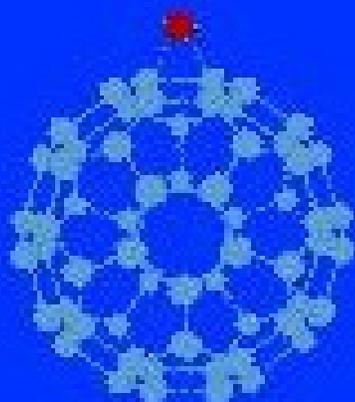
Goal: Insight into chemical phenomena.





*Exploring Chemistry with
Electronic Structure Methods*

Second Edition



*James B. Foresman
and Aileen Frisch*



**Exploring Chemistry with
Electronic Structure Methods**

Second Edition

James B. Foresman

Aileen Frisch



Gaussian, Inc.
Pittsburgh, PA

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Hamiltonian for a system with N-particles

$$\hat{H}(r)\Psi(r) = E\Psi(r)$$

Time-Independent Schrödinger Equation

$$\hat{H} = \hat{T} + \hat{V}$$

Sum of kinetic (T) and potential (V) energy

$$\hat{T} = \sum_{i=1}^N \hat{T}_i = -\sum_{i=1}^N \frac{\hbar^2}{2m_i} \nabla_i^2 = -\sum_{i=1}^N \frac{\hbar^2}{2m_i} \left(\frac{\partial^2}{\partial x_i^2} + \frac{\partial^2}{\partial y_i^2} + \frac{\partial^2}{\partial z_i^2} \right)$$

Kinetic energy

$$\nabla_i^2 = \left(\frac{\partial^2}{\partial x_i^2} + \frac{\partial^2}{\partial y_i^2} + \frac{\partial^2}{\partial z_i^2} \right)$$

Laplacian operator

$$\hat{V} = \sum_{i=1}^N \sum_{j>1}^N V_{ij} = \sum_{i=1}^N \sum_{j>1}^N \frac{q_i q_j}{r_{ij}}$$

Potential energy

$$\psi = \psi(\vec{R}_1, \vec{R}_2, \dots, \vec{R}_m, \vec{r}_1, \vec{r}_2, \dots, \vec{r}_n)$$

$$H = T_n + T_e + V_{ee} + V_{nn} + V_{en}$$

$$T_n = -\sum_{I=1}^m \frac{1}{2M_I} \nabla_I^2$$

$$T_e = -\sum_{i=1}^n \frac{1}{2} \nabla_i^2$$

$$V_{ee} = \sum_{i=1}^n \sum_{j<i} \frac{1}{|\vec{r}_i - \vec{r}_j|}$$

$$V_{nn} = \sum_{I=1}^m \sum_{J<I} \frac{Z_I Z_J}{|\vec{R}_I - \vec{R}_J|}$$

$$V_{en} = -\sum_{i=1}^n \sum_{I=1}^m \frac{1}{|\vec{r}_i - \vec{R}_I|} = -\frac{1}{|\vec{r}_i - \vec{R}_i|} - \frac{1}{|\vec{r}_J - \vec{R}_J|}$$

Born-Oppenheimer Approximation (1927)

- So far, the Hamiltonian contains the following terms:

$$\hat{H} = \hat{T}_n + \hat{T}_e + \hat{V}_{ne} + \hat{V}_{ee} + \hat{V}_{nn}$$

- Since nuclei are much heavier than electrons ($m_n/m_e=1836$), their velocities are much smaller.
- To a good approximation (**error < 1/1000***), the Schrödinger equation can be separated into two parts:
 - One part describes the electronic wavefunction for a fixed nuclear geometry.
 - The second describes the nuclear wavefunction, where the electronic energy plays the role of a potential energy.

Born-Oppenheimer Approx. cont.

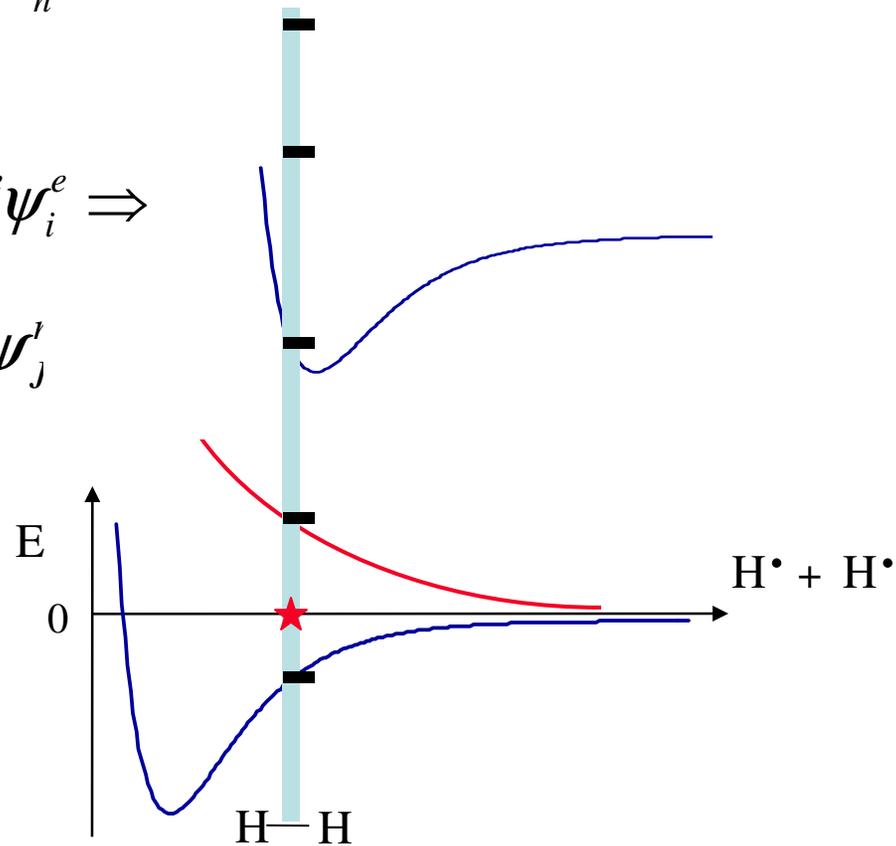
$$H\psi = E\psi$$

$$H = H_e + T_n$$

$$\psi = \psi_i^e \psi_j^n$$

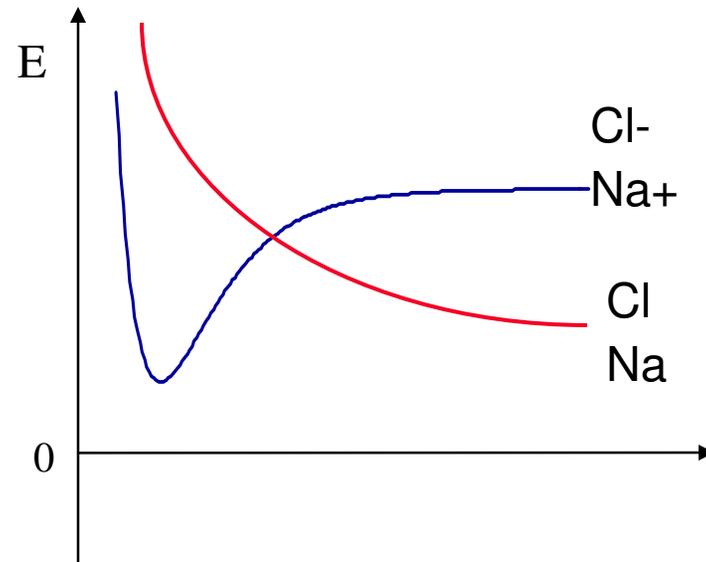
$$e: H_e \psi_i^e = E_i^e \psi_i^e \Rightarrow$$

$$n: (T_n + E_i^e) \psi_j^n$$



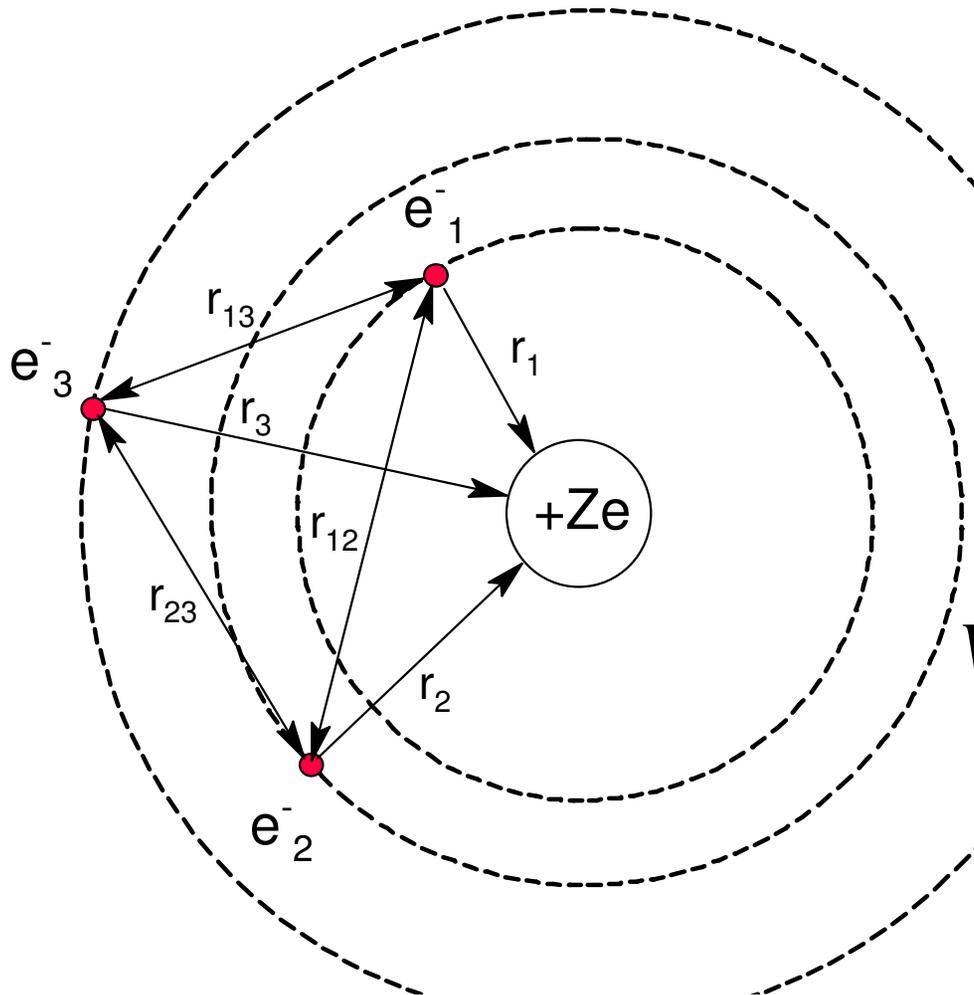
Limitations of the Born-Oppenheimer approximation

- The BO approx. is usually very good, but breaks down when two (or more) electronic states are close in energy at particular nuclear geometries. In such situations, a “non-adiabatic” wavefunction - a product of nuclear and electronic wavefunctions - must be used.



Hartree-Fock Theory

Electron-electron interaction:



$$V(r_1, r_i) = \frac{e^2}{4\pi\epsilon_0} \left[-\frac{Z}{r_1} + \sum_{i+1} \frac{1}{r_{1i}} \right]$$

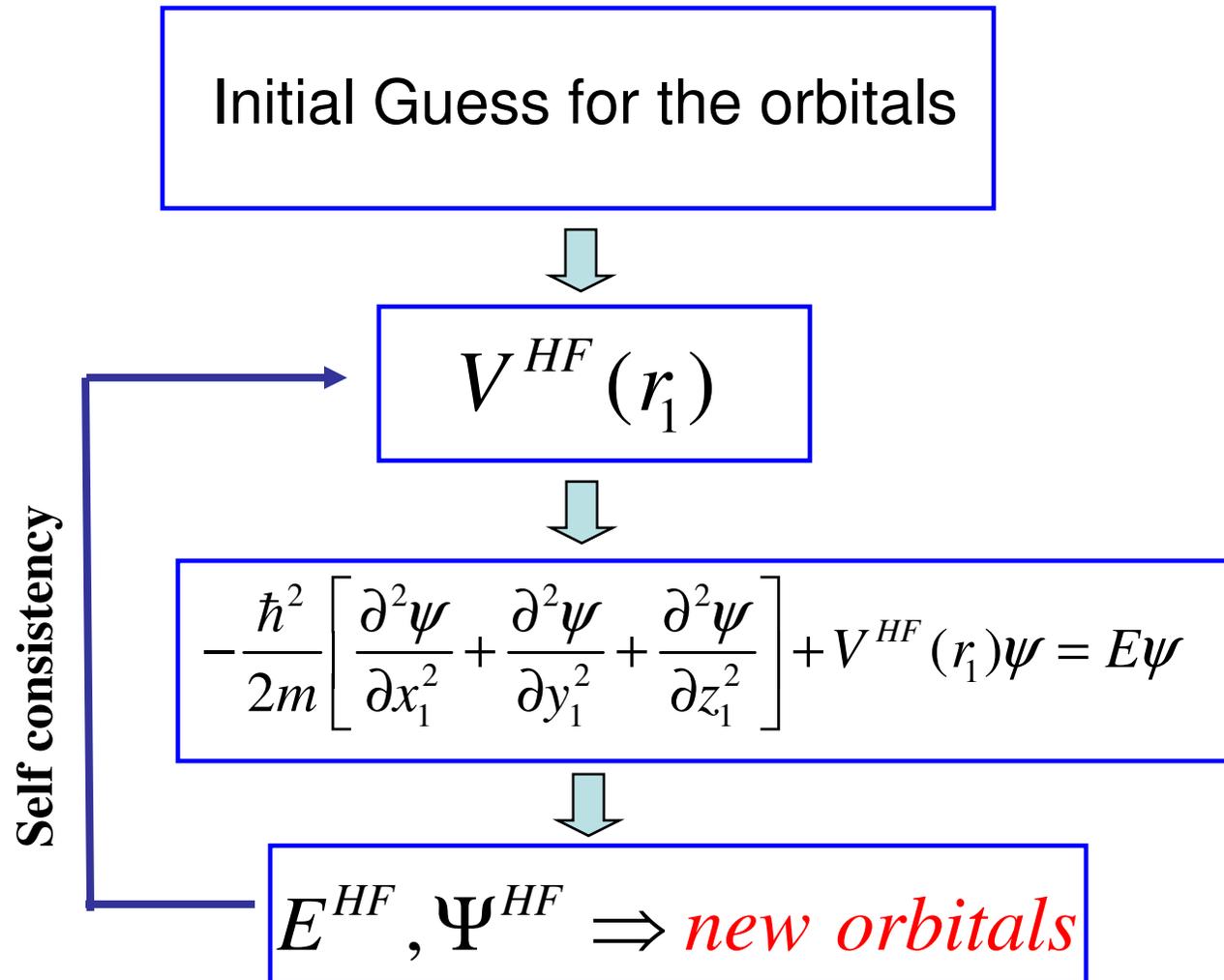


$$V(r_1) \approx \frac{e^2}{4\pi\epsilon_0} \left[-\frac{Z}{r_1} + \overline{\sum_{i+1} \frac{1}{r_{1i}}} \right]$$



$$V^{HF}(r_1) \approx \frac{e^2}{4\pi\epsilon_0} \left[-\frac{Z}{r_1} + \overline{\sum_{i+1} \frac{1}{r_{1i}}} \right]$$

Self-Consistent-Field



Restricted and Unrestricted Hartree-Fock

Restricted Hartree-Fock (RHF)

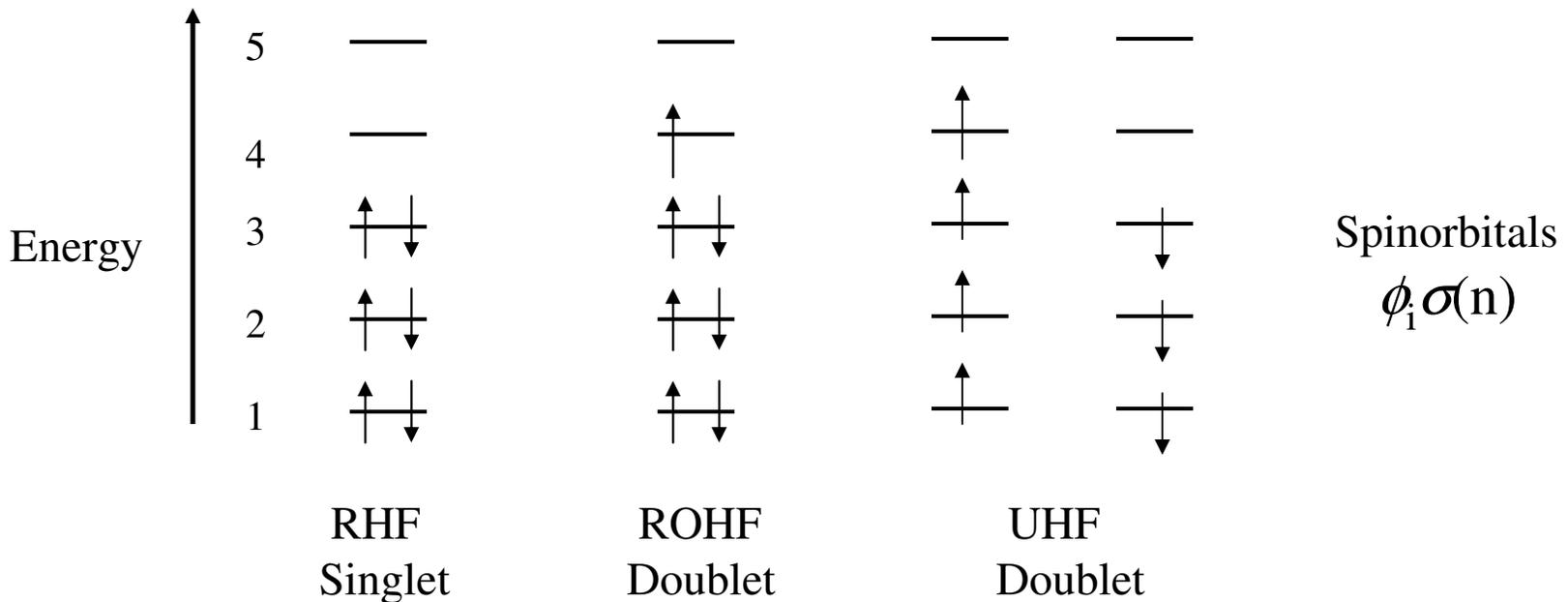
For even electron, closed-shell singlet states, electrons in a given MO with α and β spin are constrained to have the same spatial dependence.

Restricted Open-shell Hartree-Fock (ROHF)

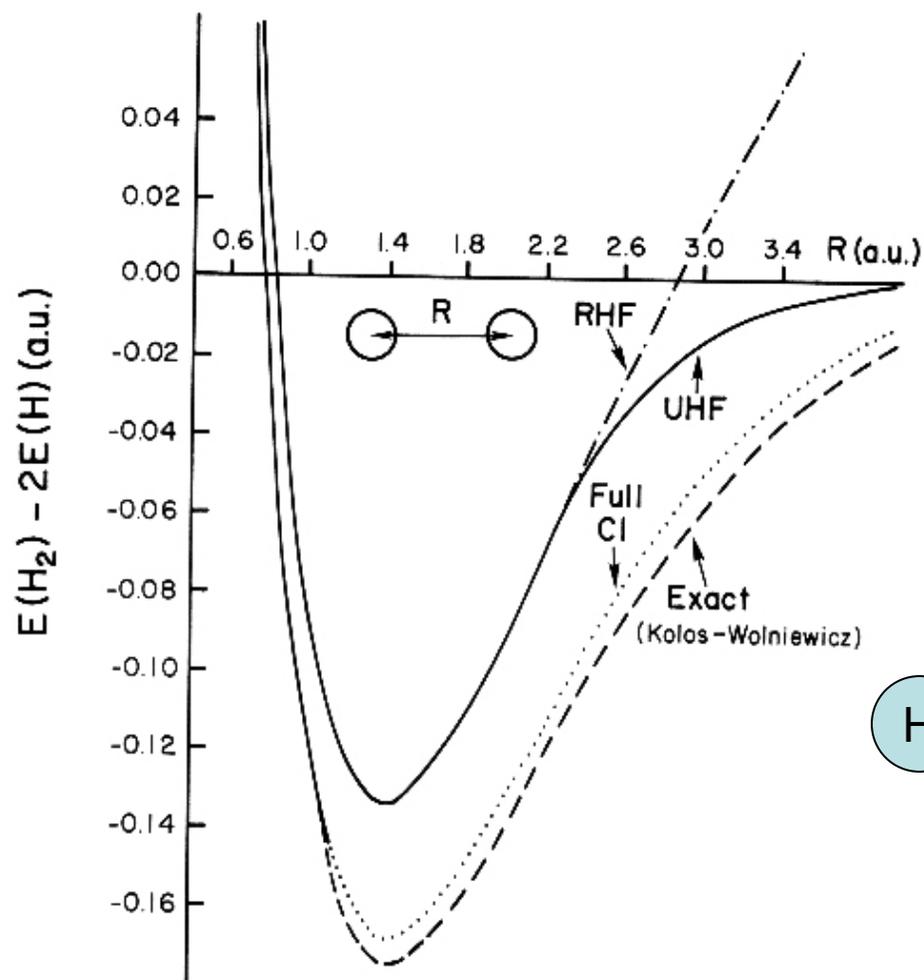
The spatial part of the doubly occupied orbitals are restricted to be the same.

Unrestricted Hartree-fock (UHF)

α and β spinorbitals have different spatial parts. α β



Potential energy curves



Restricted HF does not 'dissociate' correctly

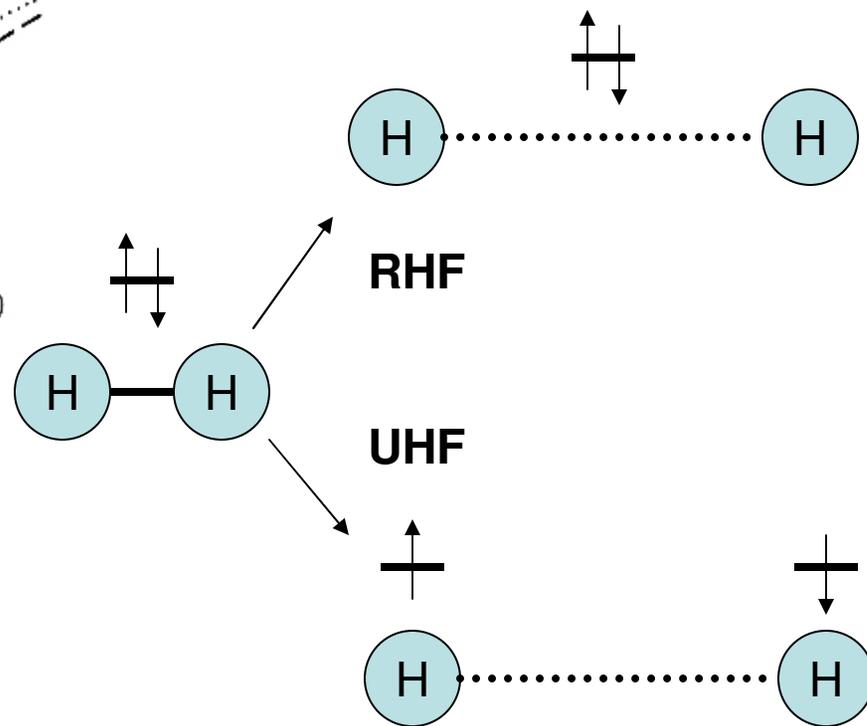


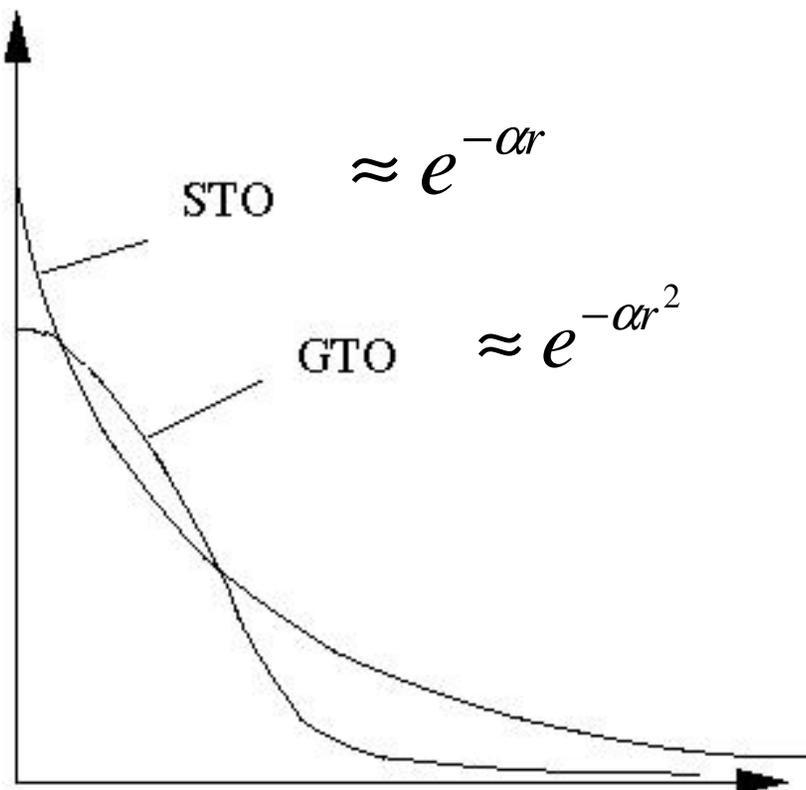
Figure 4.3 6-31G** potential energy curves for H₂.

Basis sets

$$\psi(x, y, z) = \sum_n c_n \Phi_n(x, y, z)$$

- **Mathematical trick** : Instead of looking for an unknown functions (ψ) we are searching for coefficients (c).
- “**Atoms-in-Molecule**” approach: Atomic electron density is only weakly perturbed in formation of molecule
- Hydrogen-like functions centered at atoms in molecule
- Additional functions for specific purposes (polarization, correlation, bonding etc.)

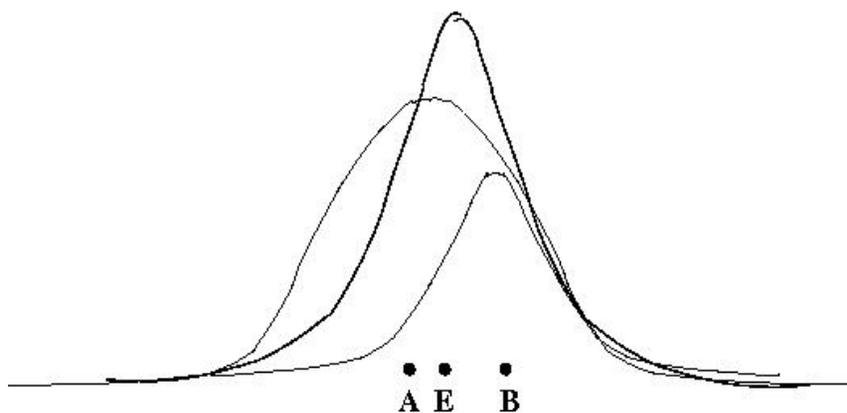
Basis sets: STOs and GTOs



- Slater-Type Orbitals are H like
- Gaussian Type Orbitals
- α : 'exponent'

Basis sets: Why GTOs ?

Situation in molecules:



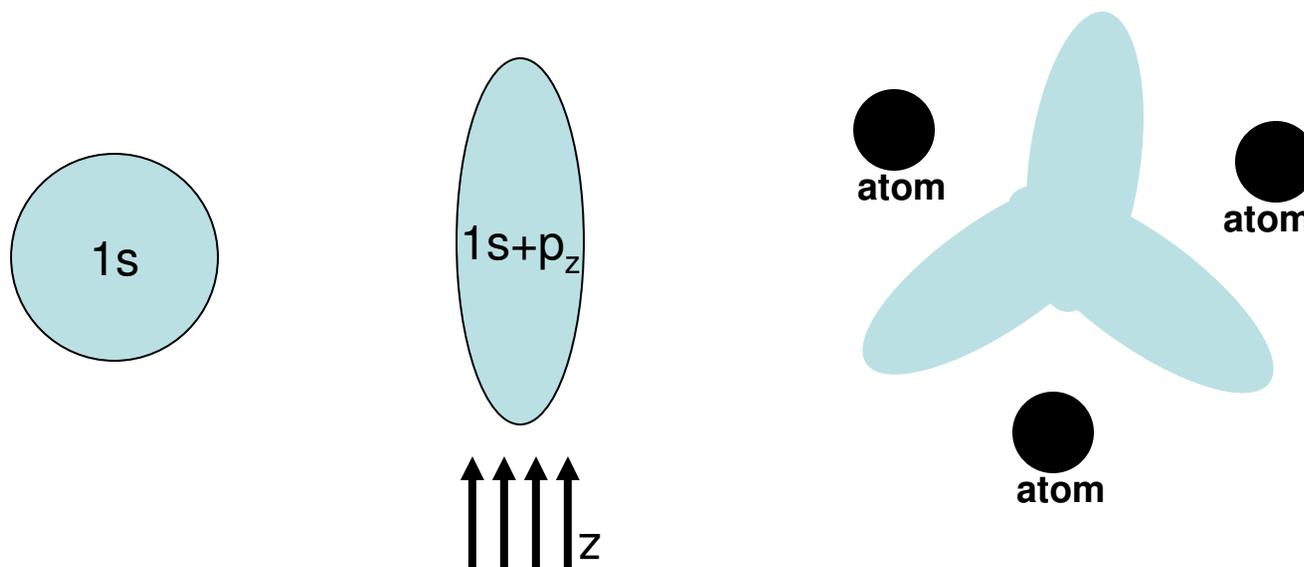
Atomic centers A,B

- Use of atom-centered GTOs
- Products of two GTOs at centers A, B gives GTO at center E
- Hence: Ease of integration!
- Much faster than with STOs, despite increased number

Standard basis sets: Overview

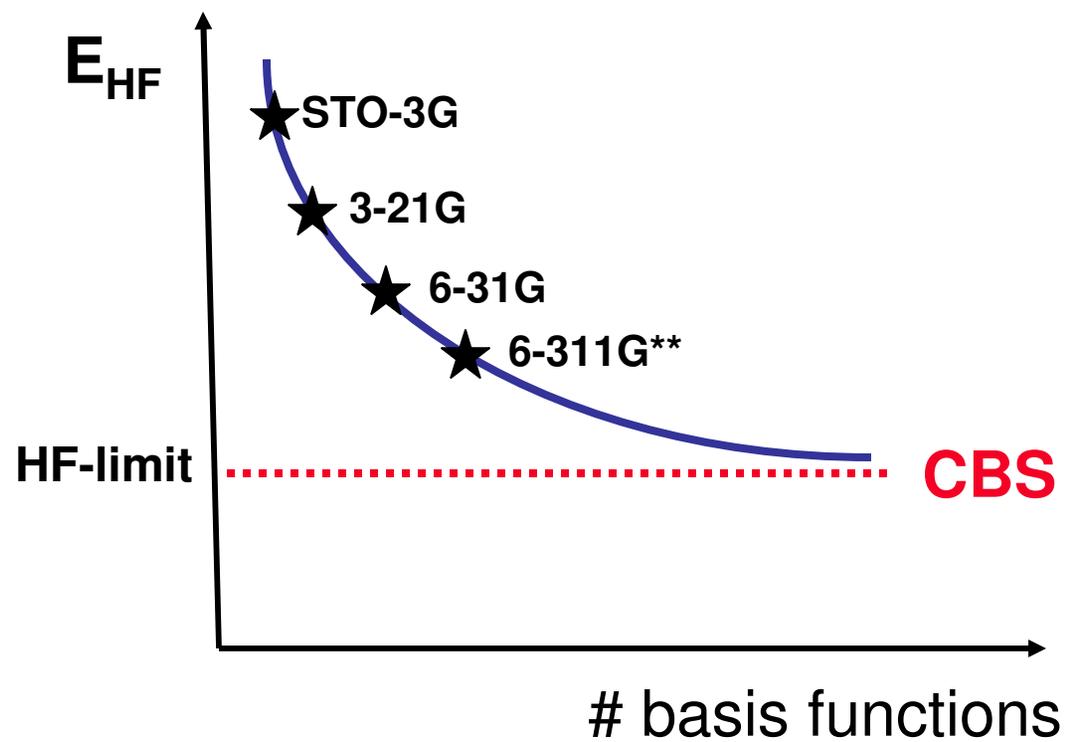
- STO-NG: N GTO per STO
- Split-valence 6-31+G(*) sets:
SZ (core)/DZ(valence),
+: diffuse functions
*: polarization functions
Z(zeta): Number of contracted functions
- (aug)-cc-(p)VXZ: Correlation-consistent basis sets

Basis sets: Polarization functions



- Polarization of atomic density upon formation of chemical bond
- Molecular field breaks atomic symmetry !
- Use of higher angular-momentum functions
E.g. p-functions for H , d-functions for C, etc.
- **Allows orbitals to change shape**

Complete Basis Set limit



Configuration Interaction (CI)

- Principles:

Based on HF (or MCSCF) wave function (orbitals)

Linear expansion of many-particle wave function:

$$\Psi_{CI} = c_0 \Psi_0 + \sum_{ar} c_a^r \Psi_a^r + \sum_{\substack{a < b \\ r < s}} c_{ab}^{rs} \Psi_{ab}^{rs} + \sum_{\substack{a < b < c \\ r < s < t}} c_{abc}^{rst} \Psi_{abc}^{rst} + \dots$$

c : expansion coefficients ($c_0 > 90\%$)

Ψ : Slater determinants, ground and excited

CI: Excitation Level

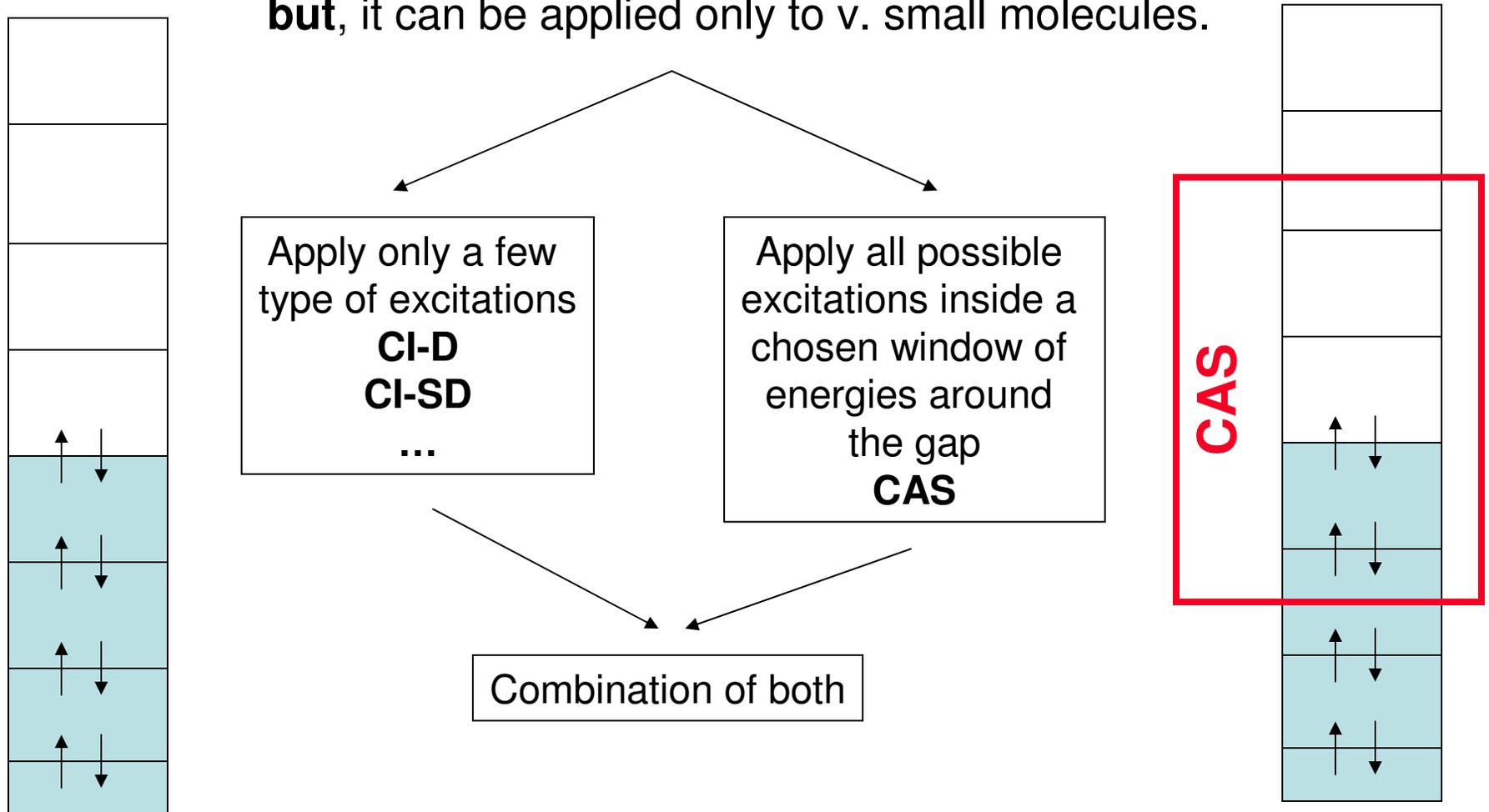
The Hamiltonian matrix in determinant basis

Anregung	$ \Psi_0\rangle$	$ S\rangle$	$ D\rangle$	$ T\rangle$	$ Q\rangle$...
$\langle\Psi_0 $	$\langle\Psi_0 \hat{H} \Psi_0\rangle$	0	$\langle\Psi_0 \hat{H} D\rangle$	0	0	...
$\langle S $		$\langle S \hat{H} S\rangle$	$\langle S \hat{H} D\rangle$	$\langle S \hat{H} T\rangle$	0	...
$\langle D $			$\langle D \hat{H} D\rangle$	$\langle D \hat{H} T\rangle$	$\langle D \hat{H} Q\rangle$...
$\langle T $				$\langle T \hat{H} T\rangle$	$\langle T \hat{H} Q\rangle$...
$\langle Q $					$\langle Q \hat{H} Q\rangle$...
\vdots	\vdots	\vdots	\vdots	\vdots	\vdots	\vdots

- Double excitations couple directly to reference state
- All others couple indirectly (coupled linear equations)
- Strategy: Select important higher excitations !

Full-CI: the computational limitation

Full CI is the **most accurate** method for a define basis set **but**, it can be applied only to v. small molecules.



Moller-Plesset Perturbation Theory

Electron correlation:
$$V = \sum_{i < j} \frac{1}{r_{ij}} - \sum_i v^{HF}$$

- Principles:

Idea: Electron correlation is a perturbation
(<1% of total energy)

Based on Rayleigh-Schrödinger PT

Single-reference method !

Size-consistent.

MPPT: Procedure

- Assume solution of zero-order problem (HF problem):

$$H_0 \Psi_0 = \sum_i f(i) \Psi_0 = E_0^{(0)} \Psi_0$$

- Definition of a perturbing potential :

$$V = \sum_{i < j} \frac{1}{r_{ij}} - \sum_i v^{HF}$$

- Total electronic Hamiltonian:

$$H = H_0 + V$$

- Apply Rayleigh-Schrödinger PT

$$E_i^{(2)} = \sum_{n \neq i} \frac{\left| \langle \Psi_i^{(0)} | \widehat{V} | \Psi_n^{(0)} \rangle \right|^2}{E_i^{(0)} - E_n^{(0)}}$$

Density Functional Theory

Hohenberg & Kohn* Theorem (1964): The ground state properties of every system are function of its charge density (ρ)

- Total energy as functional of the electron density:

$$E[\rho] = T[\rho] + U[\rho] + E_{xc}[\rho]$$

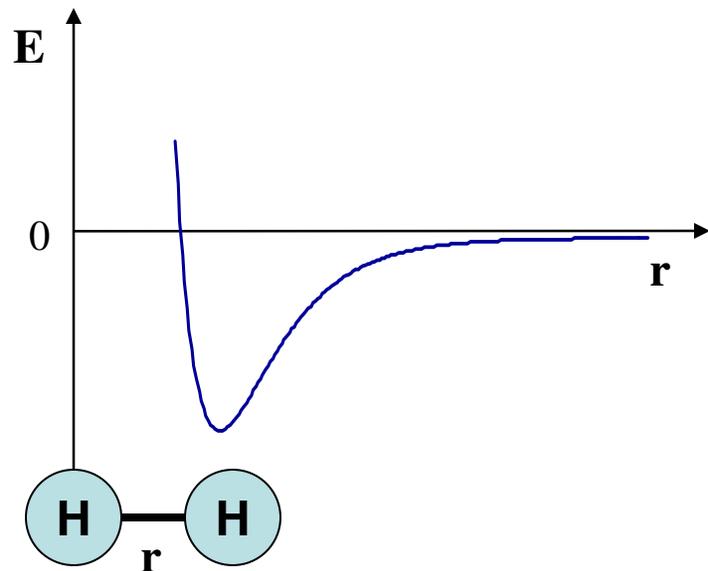
$$\rho(r) = \sum_i |\varphi_i(r)|^2 \quad : \Phi = \text{occupied MOs}$$

- T: Kinetic Energy
- U: Coulomb interactions
- xc: Exchange & correlation interactions

DFT: XC functionals

- **Local Density Approximation (LDA)**
No gradient of electron density in E_{xc}
Like a uniform electron gas !
- **Generalized Gradient Approximation (GGA)**
Gradient corrected density
- **Hybrid Functionals**
Becke (B) LYP, B3LYP functionals
Contain a contribution of HF exchange interaction

Geometry optimization



Structural & Electronic Properties

